

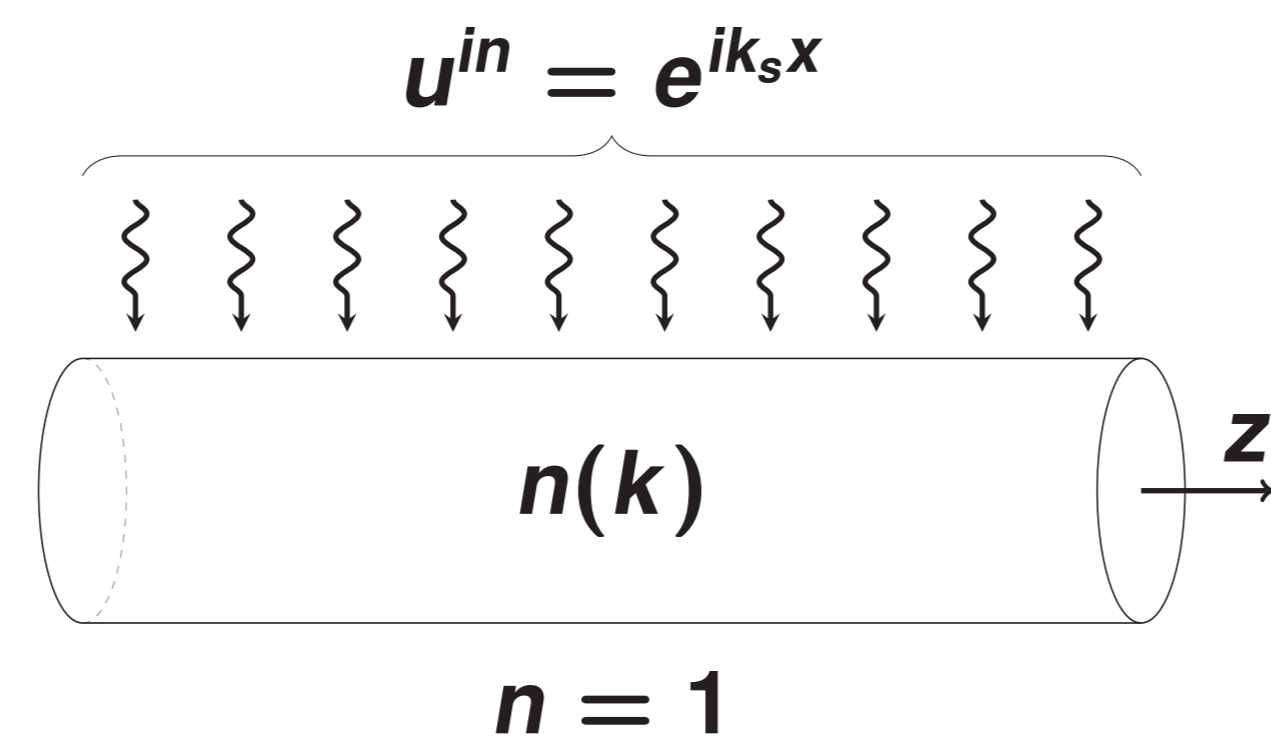
## 1) Scattering in dielectric resonators (TM)

**Objective:** Accurate computation of resonances in nano structures. Specifically, we study light confinement in prism-scatterers with arbitrarily shaped faces. Applications: design optimization of sensors, optical filters.

### Model:

- Maxwell equations in Fourier space  $\partial_t \rightarrow -i\omega$  with  $k^2 = \epsilon_0\mu_0\omega^2$
- Infinitely long dielectric ( $\partial_z \rightarrow 0$ ) in exterior domain
- TM-polarization  $\mathbf{E} = [0, 0, u(x, y)]$ ,  $H_z = 0$ . We reduce to 2D, where  $u$  satisfies:  $\mathcal{L}(k)u := \Delta u + k^2 n^2(r, k)u = 0$

- $n(r, k)$ : refractive index, where  $n \neq 1$ , for  $r \in \Omega_s$  (dielectric)
- Scattering problem:  $k_s \in \mathbb{R}$  and  $u^i$  are given, then we solve for **outgoing** solutions  $u^{out}$  of  $\mathcal{L}(k_s)u = 0$



## 3) Motivation: Energy peaks

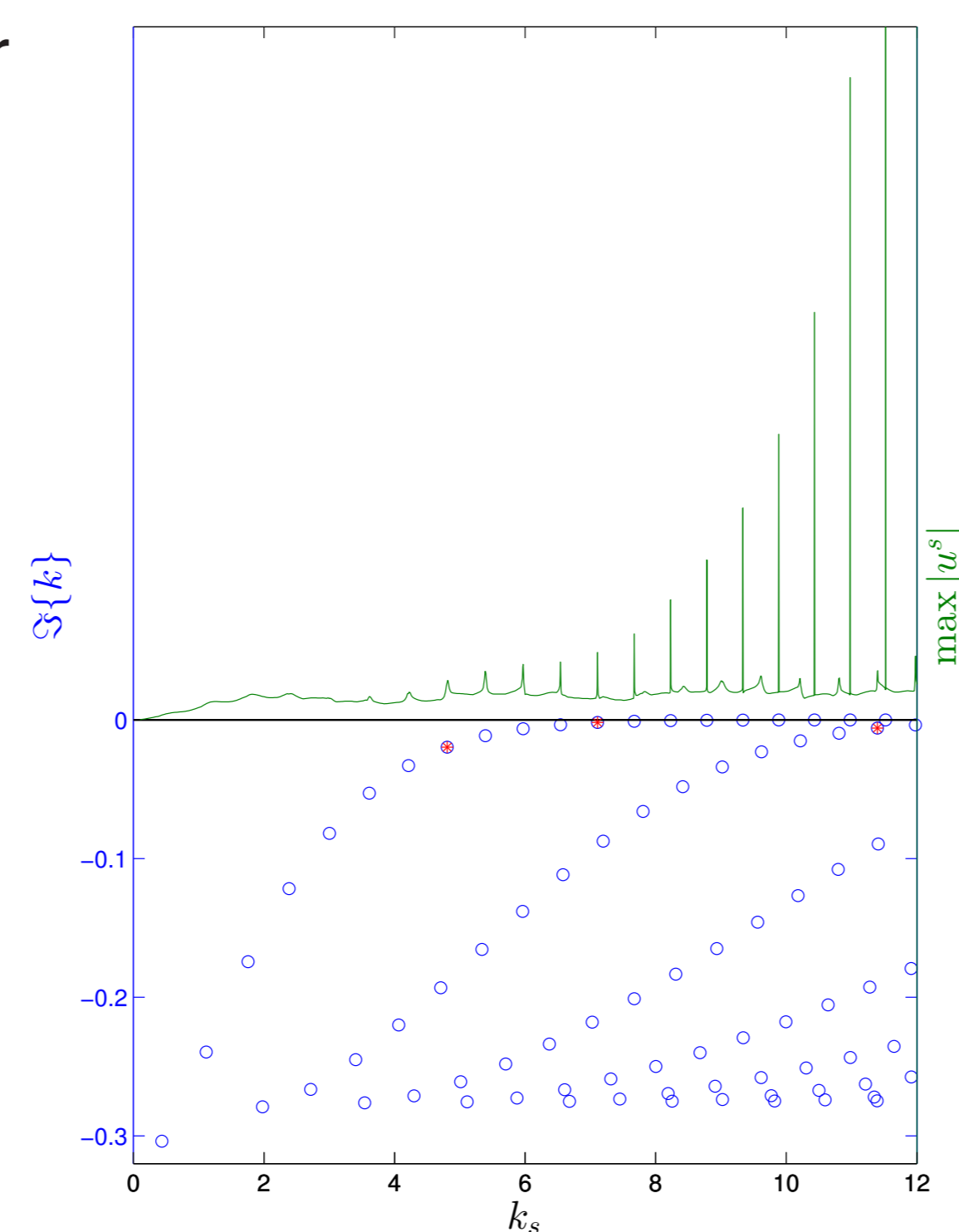
- Incident plane waves with wave number  $k_s \in \mathbb{R}$  excite amplitude peaks.
- Finding  $k_s$  by frequency sweeps is generally computationally expensive.

- Resonator's quality factor:

$$Q_m = \frac{\Re\{k_m\}}{2|\Im\{k_m\}|}$$

- Peaks and resonances  $k_m$  are related by the scattering problem through  $Q_m$ :

Take  $k_s = \Re\{k_m\}$  provided  $Q_m$  is large.

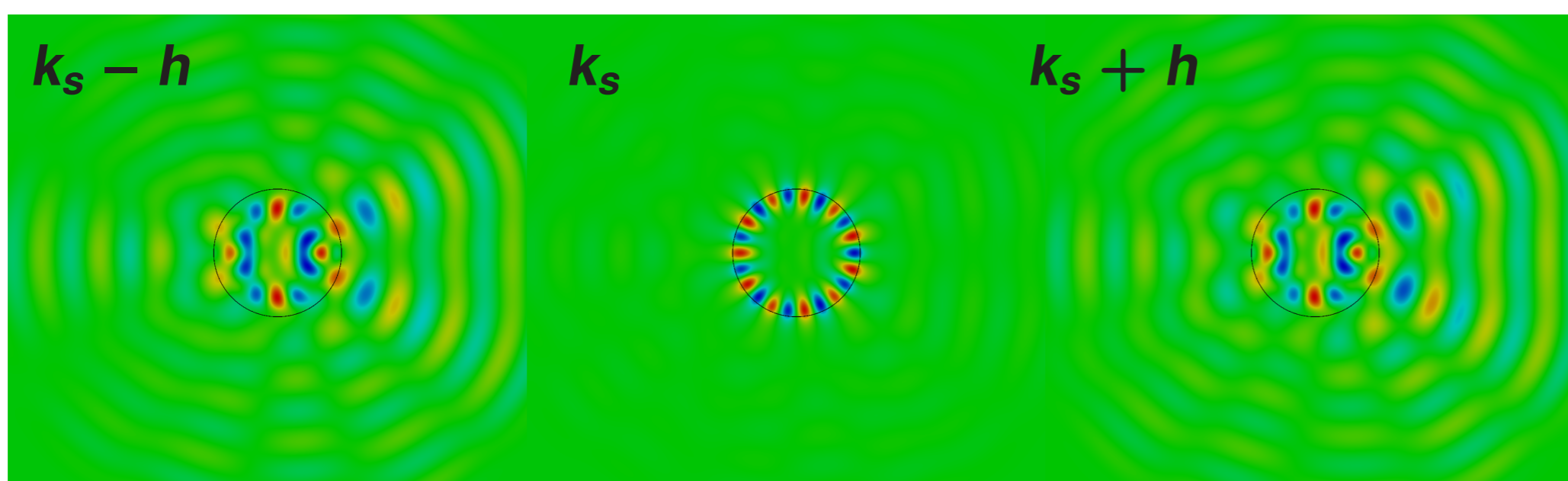


Scattering example

$$k_m = 7.109 - 0.0018i$$

$$h = 0.005$$

$$k_s = 7.109$$



## 5) Modeling exterior domains: Dirichlet-to-Neumann map

- Introduce an outer circle  $\Gamma_a$  with radius  $r = a$  containing the scatterer  $\Omega_s$  and use the exact series for outgoing waves in  $r \in \Omega^+ \equiv \mathbb{R}^2 \setminus \Omega_a$

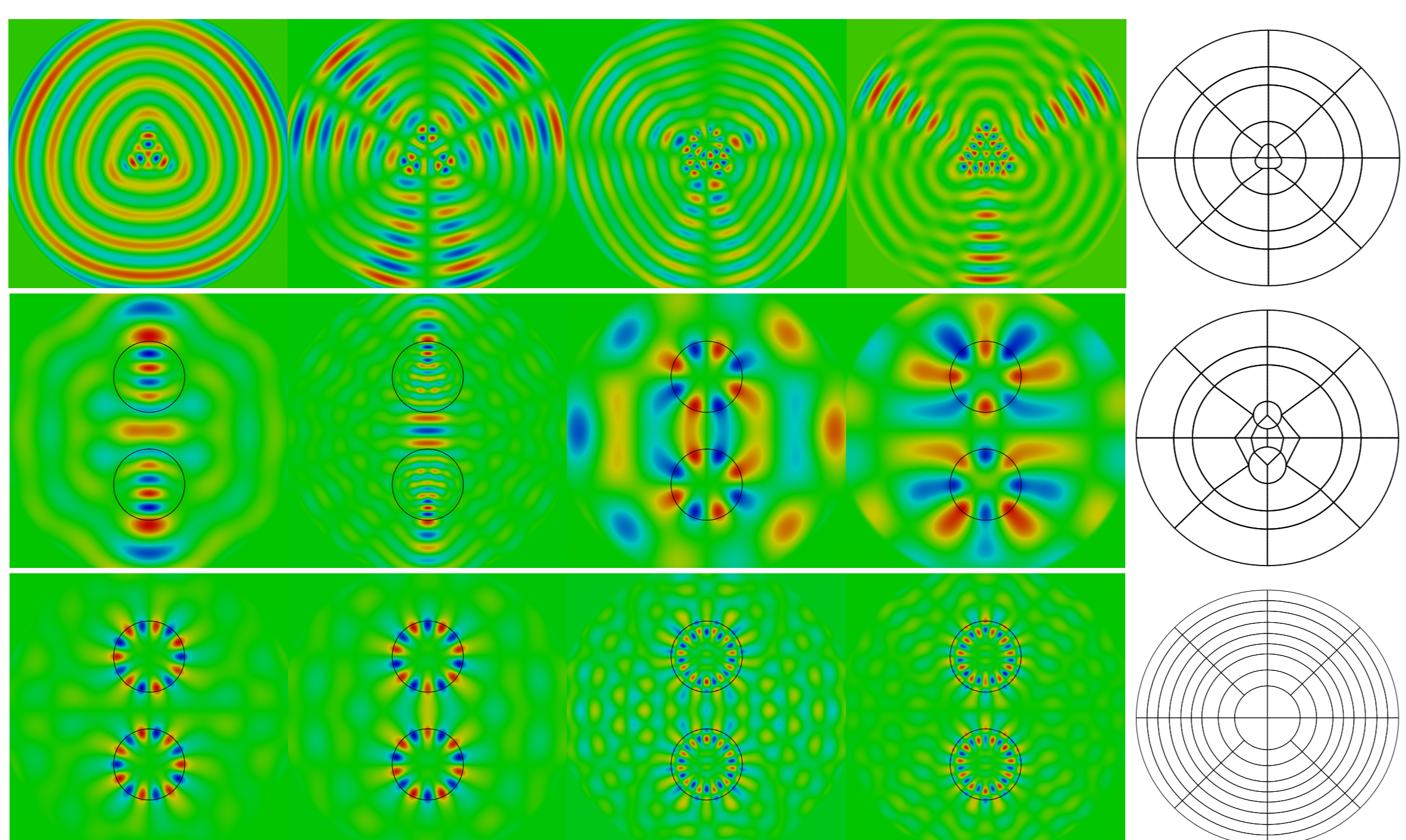
$$\Delta u + k^2 n^2(k, r)u = 0, \quad r \in \Omega_a$$

$$\nabla u \cdot \mathbf{n} = \mathcal{G}u, \quad r \in \Gamma_a \quad (\text{DtN map})$$

$$\mathcal{G}u \equiv -\frac{\partial u^+}{\partial r} \Big|_{r=a} = -\frac{1}{2\pi} \sum_{m=-\infty}^{\infty} k \frac{H'_m(ka)}{H_m(ka)} e^{im\theta} \int_0^{2\pi} u(a, \theta') e^{-im\theta'} d\theta'$$

- Truncate the series until  $|m| = l$ , with  $l > a\Re\{k\}$
- Weak form: Find eigenpairs  $(u, k)$  such that: (nonlinear in  $k$ )  $(\nabla u, \nabla \bar{v})_{\Omega_a} - k^2(n^2(k, r)u, \bar{v})_{\Omega_a} + (\mathcal{G}^l(k)u, \bar{v})_{\Gamma_a} = 0, \quad u, v \in H^1(\Omega_a)$
- Nonlinear eigenvalue problem, solver: NEPCISS in SLEPc-3.6.0.

## 7) Resonances of arbitrarily shaped scatterers



- We use curved elements (*Manifolds* in Deal.II library) and polynomial order  $p = 20$

## 2) Problem description and challenges

- A **resonance** state is a long-lived state of an open system and can be determined by applying a DtN-map formulation or by a perfectly matched layer (PML) to  $\mathcal{L}(k)u = 0$ ; resonances  $k_m \in \mathbb{C}$ .
- Solutions  $u_m \notin L^2(\mathbb{R}^2)$  grow as  $e^{|\Im\{k_m\}|d}$ , with  $d$  the distance from  $\Omega_s$ . This makes the  $u_m$  difficult to compute by any discretization method.
- Sectors in the  $k$ -plane may have large resolvent norms, which results in spurious solutions for poor discretizations (spectral pollution). Example: allowing large *air* regions in the computational domain.
- Identifying spurious solutions is an unsolved problem.
- We discretize by using FEM with high polynomial degree. We use *Gauss-Lobatto* shape functions and curvilinear elements in **deal.II**.

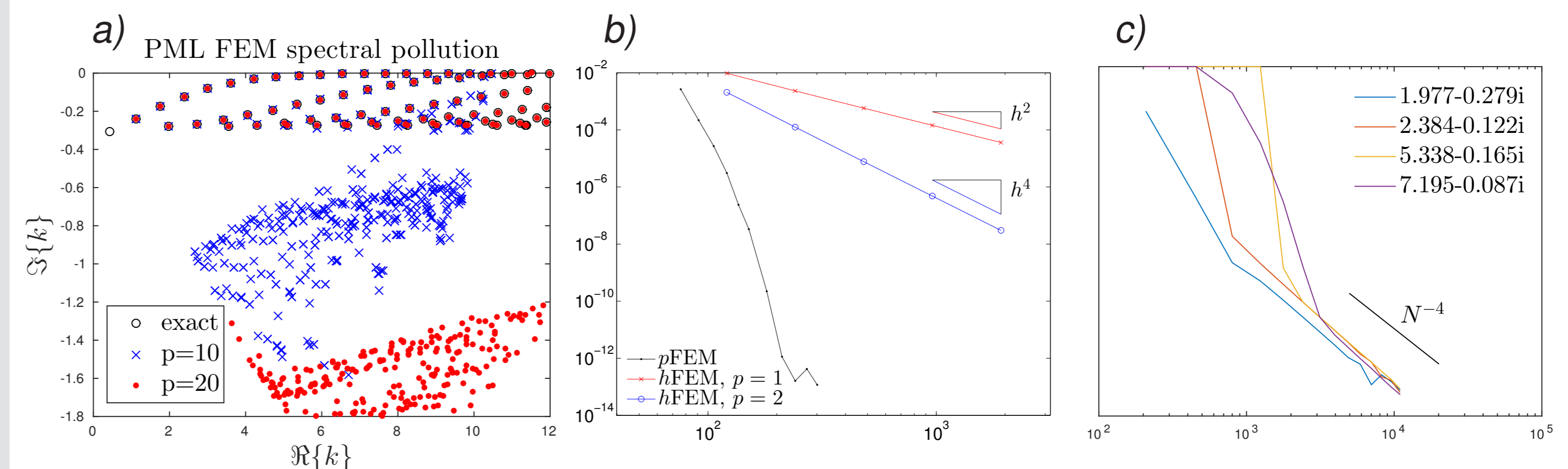
## 4) Modeling exterior domains: Radial PML

- Coordinate stretching  $\nabla \cdot (\mathcal{A}\nabla u) + k^2 n^2 \mathcal{B}u = 0, \quad r \in \Omega_d$   
with definitions  $\alpha := 1 + i\sigma(r), \quad u = 0, \quad r \in \Gamma_d$

$$\mathcal{A}(r) = \begin{pmatrix} \frac{\tilde{\alpha}}{\alpha} \cos^2 \theta + \frac{\alpha}{\tilde{\alpha}} \sin^2 \theta & (\frac{\tilde{\alpha}}{\alpha} - \frac{\alpha}{\tilde{\alpha}}) \sin \theta \cos \theta \\ (\frac{\tilde{\alpha}}{\alpha} - \frac{\alpha}{\tilde{\alpha}}) \sin \theta \cos \theta & \frac{\tilde{\alpha}}{\alpha} \sin^2 \theta + \frac{\alpha}{\tilde{\alpha}} \cos^2 \theta \end{pmatrix}, \quad \mathcal{B}(r) = \alpha \tilde{\alpha}$$

- Weak form: Find eigenpairs  $(u, k^2)$  s.t:  
 $(\mathcal{A}\nabla u, \nabla \bar{v})_{\Omega_d} - k^2(n^2(k, r)\mathcal{B}u, \bar{v})_{\Omega_d} = 0; \quad u, v \in H_0^1(\Omega_d)$
- $\mathcal{A}$  is the identity in air/scatterer, and has diagonal form for  $r > r_2$
- Parameters: domain's truncation  $d$  and strength  $\sigma_0$
- Generalized linear eigenvalue problem for  $n(k, r) \equiv n(r)$  and  $\lambda \equiv k^2$
- Solver: *EPSKRYLOV*SCHUR in SLEPc.

## 6) Convergence and comparison



- In a) we show examples of spectral pollution from PML computations. It is evident that improving the mesh quality (increasing the polynomial degree  $p$ ) moves spurious eigenvalues away from the area of interest.
- In b) we show relative errors for the 3<sup>rd</sup> eigenvalue in the PML 1D case. The simplicity of the problem allow us to study convergence for i) *hFEM*: where we keep  $p$  fixed and decrease the mesh size  $h$ , and ii) *pFEM*: where  $h$  is fixed and we increase  $p$ . In i) we obtain  $h^{2p}$  convergence, where in ii) we obtain exponential convergence.
- In c) we show relative errors for the 2D disk computed with the DtN method and  $l = 30$  (truncation index) for selected eigenvalues. We used the nonlinear eigenvalue solver NEPCISS in SLEPc.

## 8) PML vs DtN for linear materials: $n(r, k) \equiv n(r)$

- PML:**
- + Linear eigenvalue problem.
  - + Easy to implement.
  - Nonphysical eigenvalues satisfy the PML PDE.
  - Resonances with small  $|\Re\{k\}|$  require large  $\sigma_0$ , more dofs!
  - Parameters  $\sigma_0$  and  $d$  are not straight forward to control.
  - Requires extra cells/dofs for the PML layer.
  - Spurious eigenvalues appear for poor discretizations.
- DtN:**
- + Exact boundary condition, only resonances are eigenvalues.
  - + The only parameter is  $l$  (truncation of the series).
  - + Only physical cells are required (fewer dofs).
  - Resonance problem: nonlinear eigenvalue problem (quadratic in 1D).
  - Implementation requires keeping  $2l + 1$  DtN dense boundary matrices.
  - Spurious eigenvalues appear for poor discretizations.
- Spectral pollution can be diminished by using finer discretizations.
  - Eigenvalue computation of non-self adjoint operators is a challenging task.